

# UNIVERSITY OF FORT HARE



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**PHY 502**

## HONOURS EXAMINATION

DATE : JUNE/JULY 2023  
TIME : 3 HOURS  
SUBJECT : PHYSICS 502 Electrodynamics  
MARKS : 100

INTERNAL EXAMINERS:

Prof G Makaka

EXTERNAL EXAMINER

Dr D Tinarwo

**INSTRUCTIONS:** Answer all the four question, each question carries 25 marks

1. (a) Derive the general inhomogeneous wave equations for the scalar and vector potentials in electrodynamics. [13]
- (b) Then, specialise to the Lorentz condition and write down the wave equations for this particular condition. [7]
- (c) Finally, gauge transform the Lorentz equations just derived. [5]
2. (a) State what is giving rise to the current density  $\vec{j}$ , which is given at any point  $\vec{r}'$  and time  $t$ , in the usual notation, by  $\vec{j}(\vec{r}', t) = q\dot{\vec{r}}(t)\delta(\vec{r}' - \vec{r}(t))$ . [4]
- (b) Write the expression for the four-current density  $j^\mu$  in terms of charge density  $\rho$  and current density  $\vec{j}$ . [4]
- (c) Write the charge conservation equation in covariant form. [4]
- (d) In a certain region of space, the electric scalar potential and the magnetic vector potential are given as:  $\Phi(\vec{r}, t) = \frac{\Phi_0 e^{i(\omega t - kr)}}{kr}$  and  $A(\vec{r}, t) = \frac{\vec{A}_0 e^{i(\omega t - kr)}}{kr}$ , where  $\Phi_0$  and  $\vec{A}_0$  are both constant and  $r = |\vec{r}|$ . Find corresponding expressions for the  $\vec{E}$  and  $\vec{B}$  fields. [13]

3. (a) (i) Show that the tensor scalar product of the electromagnetic field tensor is:

$$F_{\alpha\beta} F^{\beta\alpha} = 2 \left( \frac{E^2}{c^2} - B^2 \right), \quad [5]$$

- (ii) Hence show that  $\tilde{F}_{\alpha\beta} \tilde{F}^{\beta\alpha} = -F_{\alpha\beta} F^{\beta\alpha}$  where  $\tilde{F}_{\alpha\beta}$  is the dual of  $F_{\alpha\beta}$ . [5]

- (b) Show That Maxwell's homogeneous equations can be written in the form:

$$\partial^\alpha F^{\beta\gamma} + \partial^\beta F^{\gamma\alpha} + \partial^\gamma F^{\alpha\beta} = 0 \quad [7]$$

- (c) Show that  $\partial_\alpha \tilde{F}^{\alpha\beta} = \begin{bmatrix} c\nabla \cdot \vec{B} \\ -(\nabla \times \vec{E})_x - \partial_t B_x \\ -(\nabla \times \vec{E})_y - \partial_t B_y \\ -(\nabla \times \vec{E})_z - \partial_t B_z \end{bmatrix}$  [8]

4. (a) The Poynting's theorem can be written in the form  $\frac{dW}{dt} = -\frac{d}{dt} \int_V U_{em} d\tau - \int_{surface} \vec{s} \cdot d\vec{a}$ , show

that it can be expressed in the form:  $\nabla \cdot \vec{S} = -\partial_t U$ , where  $U = U_{mech} + U_{em}$  [7]

(b) Starting from the Lorentz force law, show that the Lorentz force density  $\vec{f}$  can be expressed in the form:  $f_k = \partial_j \vec{T}_{jk} - \epsilon_0 \mu_0 \frac{\partial \vec{S}_k}{\partial t}$ , where;

$$\partial_j \vec{T}_{jk} = \epsilon_0 \left\{ (\nabla \cdot \vec{E}) \vec{E} + (\vec{E} \cdot \nabla) \vec{E} - \frac{1}{2} \nabla E^2 \right\}_k + \frac{1}{\mu_0} \left\{ (\nabla \cdot \vec{B}) \vec{B} + (\vec{B} \cdot \nabla) \vec{B} - \frac{1}{2} \nabla B^2 \right\}_k$$

[10]

(c)  $\Sigma$  and  $\Sigma'$  are space-time inertial coordinate frames such that at time  $t = t' = 0$ , the origins coincide. The electromagnetic field tensor  $F^{\mu\nu}$  has the following components in  $\Sigma$ :

$$F^{\mu\nu} = \begin{bmatrix} 0 & \frac{-E_x}{c} & \frac{-E_y}{c} & \frac{-E_z}{c} \\ \frac{E_x}{c} & 0 & -B_z & B_y \\ \frac{E_y}{c} & B_z & 0 & -B_x \\ \frac{E_z}{c} & -B_y & B_x & 0 \end{bmatrix}$$

(i) Write the covariant equation that links  $F^{\mu\nu}$  to its sources. [4]

(ii) Consider the four-vector  $(f^0, \vec{f})$  defined by  $f^\mu = F^{\mu\nu} j_\nu$ . Deduce, in a clear explicit calculation, that  $\vec{f}$  is the Lorentz force density form. [4]

END OF PAPER